

John Schwarz
Aspen 2/05

STRING THEORY - THE EARLY YEARS

A PERSONAL PERSPECTIVE

Some of the topics have been written up in three recent preprints:

hep-th/0007117 (Scherk)

hep-th/0007118 (Susskind)

hep-th/0011078 (Susskind)

I will discuss the period 1960-85 at a rate of about 2 minutes per year.

DISCLAIMER: I make no pretense of balance, fairness, or objectivity.

OUTLINE

- 1960 - 68 The analytic S matrix
- 1968 - 70 The dual resonance model
- 1971 - 73 The NSR model
- 1974 - 75 Gravity + unification
- 1975 - 79 Supersymmetry + supergravity
- 1979 - 85 Superstrings
- Concluding comments

1960-68 THE ANALYTIC S MATRIX

GOAL: Theory of hadrons

UC Berkeley was center of universe

(Gell-Mann, Mandelstam, Weinberg, Glashow...)

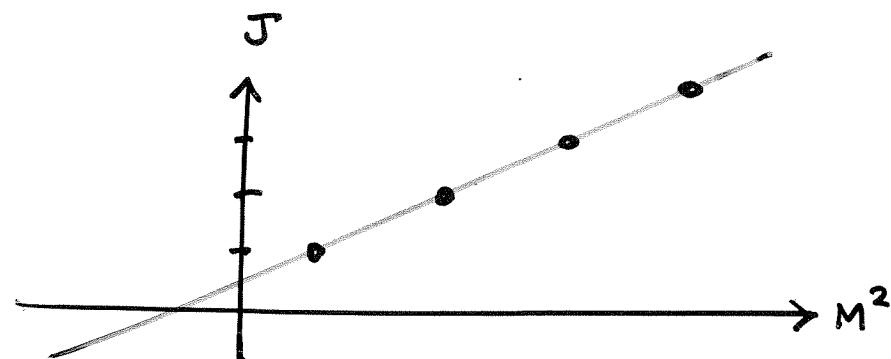
I was grad student there 1962-66

(so was David Gross)

PRINCIPLES:

- Only on-shell S Matrix is physical.
Field theory is misguided.
- Unitarity and analyticity
- Bootstrap conjecture
- Analyticity in angular momentum

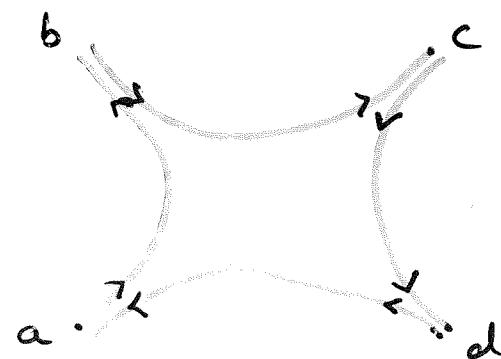
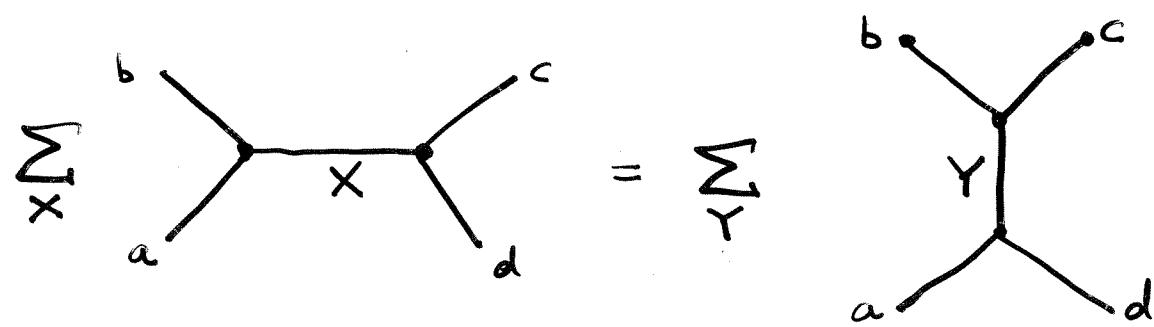
Regge Pole Theory



$$A(s, t) \sim \rho(t) (s/s_0)^{\alpha'/t}, \quad t < 0$$

$$\alpha' \sim 1.0 \text{ (GeV)}^{-2}$$

Duality



meson - meson scattering.

1968-70 THE DUAL RESONANCE MODEL

Veneziano formula

$$T = A(s, t) + A(s, u) + A(t, u)$$

$$A(s, t) = \frac{\Gamma(-\alpha(s)) \Gamma(-\alpha(t))}{\Gamma(-\alpha(s) - \alpha(t))}$$

$$\text{where } \alpha(s) = \alpha(0) + \alpha's$$

This function gives an explicit realization of duality and Regge behavior in the "narrow resonance approximation". The motivation was phenomenological, but this turned out to be a tree amplitude in a theory!

Virasoro formula

$$T = \frac{T(-\frac{1}{2}\alpha(s)) T(-\frac{1}{2}\alpha(t)) T(-\frac{1}{2}\alpha(u))}{T(-\frac{1}{2}\alpha(t) - \frac{1}{2}\alpha(u)) T(-\frac{1}{2}\alpha(s) - \frac{1}{2}\alpha(u)) T(-\frac{1}{2}\alpha(s) - \frac{1}{2}\alpha(t))}$$

has similar virtues.

- N-particle generalization of Virasoro formula:

$$A_N = \int \pi dy_i \mu(y) \prod_{i < j} |y_i - y_j|^{\alpha' k_i \cdot k_j}$$

This has cyclic symmetry in the N external lines.

- N-particle generalization of Virasoro formula

$$T_N = \int \pi d^2 z_i \mu(z) \prod_{i < j} |z_i - z_j|^{\alpha' k_i \cdot k_j}$$

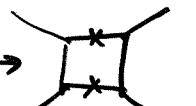
This has total symmetry in the N external lines.

Both of these formulas were shown to have a consistent factorization on a spectrum of single-particle states described by an ∞ number of oscillators

$$\{a_m^\mu\} \quad \begin{matrix} \mu = 0, 1, \dots, d-1 \\ m = 1, 2, \dots \end{matrix} \quad (2 \text{ sets for S-V case})$$

Eventually these formulas were interpreted as describing the scattering of modes of a relativistic string: open strings in the first case + closed strings in the second case. Amazingly, the formulas preceded the interpretation!

Having found the factorization,
it became possible to study
radiative corrections (loop amplitudes)

I was at Princeton & collaborated
with Gross, Neveu, & Sclar in
studying one-loop amplitudes.
In particular, we found new
singularities in the "nonplanar"
open-string loop $s \rightarrow$ .

These appeared to be unitarity-violating
branch points. Lovelace argued they
would be poles for $\alpha(10) = 1 + d = 26$.
Later they were interpreted as
due to closed strings.

1971-73

NSR MODEL

RAMOND - string analog of Dirac spin

$$P^\mu = \int_0^{2\pi} P^\mu(\sigma) d\sigma$$

$$\gamma^\mu = \int_0^{2\pi} \Gamma^\mu(\sigma) d\sigma$$

$$F_n = \int_0^{2\pi} e^{-inx} T \cdot P d\sigma , \quad n \in \mathbb{Z}$$

$$F_0 = \gamma \cdot p + \text{oscillator terms}$$

$(F_0 + m)|\psi\rangle = 0$ for a free fermionic string

Super-Virasoro algebra:

$$\{F_m, F_n\} = 2L_{m+n} + \frac{c}{3}m^2 \delta_{m,-n}$$

$$[L_m, F_n] = \left(\frac{m}{2} - n\right) F_{m+n}$$

$$[L_m, L_n] = (m-n)L_{m+n} + \frac{c}{12}m^3 \delta_{m,-n}$$

NS bosons

similar structure, but periodic $T^{\mu}(\sigma)$
replaced by antiperiodic $H^{\mu}(\sigma)$

$$G_r = \int_0^{2\pi} e^{-ir\sigma} H \cdot P d\sigma \quad r \in \mathbb{Z} + \frac{1}{2}$$

satisfy a similar super-Virasoro algebra.

The string world-sheet theory

$$S = \int d\sigma d\tau \left\{ \partial_{\alpha} X^{\mu} \partial^{\alpha} X_{\mu} - i \bar{\psi}^{\mu} \rho^{\alpha} \partial_{\alpha} \psi_{\mu} \right\}$$

has 2d supersymmetry

$$\delta X^{\mu} = \bar{\epsilon} \psi^{\mu}$$

$$\delta \psi^{\mu} = -i \rho^{\alpha} \epsilon \partial_{\alpha} X^{\mu}$$

discovered by Gervais & Sakita (1971).

Nelson, Charles Thorn, & I.
assembled these bosons & fermions
into a consistent interacting theory.
In 1972 I discovered the need for $d=10$.

- GGRT - light-cone gauge quantization of strings
- No-ghost theorem - Browne, Goddard, Thorn, THS
— moved to Caltech —
- Fermion-fermion scattering - THS & C.C. Wu

Other developments in this period:

- Completion & acceptance of the standard model
- The Webs-Zumino model

Understandably, string theory rapidly fell out of favor.

1974-75: GRAVITY & UNIFICATION

String theory of hadrons had problems:

- tachyons
- $d = 10$ or 26
- massless particles with $J \leq 2$.

Several years of attempts to do better were unsuccessful.

In 1974, when Toël Scherk was spending a half year at Caltech, we decided to try to interpret the spin 2 state as a graviton. (Yoneya,
too) To account for Newton's constant we estimated that

$$\alpha' \sim (1/m_p)^2$$

instead of $(1 \text{ GeV})^{-2}$ - a change by a factor of 10^{38} .

This had several advantages:

- gravity was required by the theory
- There are no UV divergences
- Extra dimensions could be a good thing
- YM gauge field also appeared naturally and could account for the other forces.

I ~~was~~ very excited and decided to dedicate my life to this.

It was puzzling to me — and still is — why it took 10 years to convince most other physicists that this is a good idea.

1975-79 SUSY + SUGRA

$N=1, D=4$ supergravity

Ferrara, Freedman, Van Nieuwenhuizen;

Deser, Zumino

$N=1, D=10$ + $N=4, D=4$ super YM

Brink, Scherk, + THS

$N=8, D=4$ supergravity

Cremmer, Julia

$N=1, D=11$ supergravity

Cremmer, Julia, Scherk

I also worked with Brink, Ramond,
+ Bell-Mann on a (not very
successful) study of superspace
supergravity.

Gliozzi, Scherk, Olive (1980)
 proposed a projection of the RNS
 spectrum that removed roughly
 half the states (including the tachyon).
 Then they did the following counting:

$$\begin{aligned}
 f_{NS}(w) &= \sum_{n=0}^{\infty} d_{NS}(n) w^n \\
 &= \frac{1}{2\sqrt{w}} \left(\prod_{m=1}^{\infty} \left(\frac{1+w^{m-\frac{1}{2}}}{1-w^m} \right)^8 - \prod_{m=1}^{\infty} \left(\frac{1-w^{m-\frac{1}{2}}}{1-w^m} \right)^8 \right)
 \end{aligned}$$

$$\begin{aligned}
 f_R(w) &= \sum_{n=0}^{\infty} d_R(n) w^n \\
 &= 8 \prod_{m=1}^{\infty} \left(\frac{1+w^m}{1-w^m} \right)^8
 \end{aligned}$$

In 1829, Jacobi discovered that

$$f_R(w) = f_{NS}(w)$$

Thus the number of bosons + fermions
 match at every mass level.

1979-85

SUPERSTRINGS

In 1979, following a year at the Ecole Normale, I spent a month in the summer at CERN. There Michael Green and I began a long + exciting collaboration.

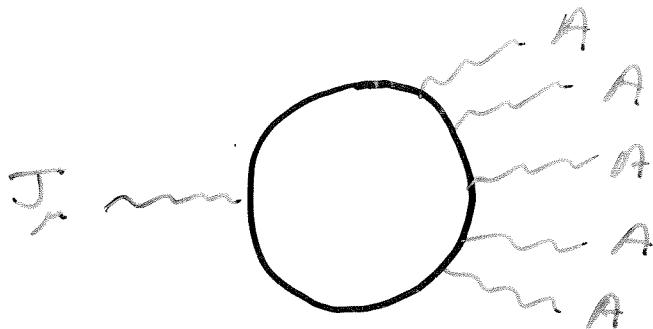
Our first goal was to understand better why the GSO projected RNS string theory has spacetime supersymmetry.

Michael had several extended stays at Caltech and I had one in London. We also worked together several summers in Aspen.

MBG + JHS (partially, also with
LARS BRINK)

- Developed a new light-cone gauge formalism in which spacetime susy of the spectrum and interaction was easily proved
- Used this formalism to compute various tree and one-loop amplitudes and elucidate their properties
- Formulated (and named) the type I, type IIA, and type IIB superstring theories
- Formulated superstring field theory in the light-cone gauge
- Formulated a covariant world-sheet theory with manifest spacetime susy.

ANOMALIES



$$\partial_\mu J^\mu \sim \epsilon^{\mu_1 \dots \mu_{10}} F_{\mu_1 \mu_2} \dots F_{\mu_9 \mu_{10}}$$

Type I superstrings are well-defined at tree-level for $G = SO(n)$ or $Sp(n)$.

In every case they are chiral (i.e., parity violating) and the $d=10$ SYM sector is anomalous.

It appeared unlikely that superstrings could accommodate parity violation, but a better analysis was required.

We worked on this problem
for almost two years, until the
crucial breakthroughs were made
in August 1984 (in Aspen).

I won't belabor the details here
— just state the conclusion:

Anomalies can cancel for a
theory with $N=1$ supersymmetry in $d=10$,
only if the YM gauge group is
 $SO(32)$ or $E_8 \times E_8$.

The $SO(32)$ case could be
accommodated by type I superstring,
but we didn't know a superstring
theory with gauge group $E_8 \times E_8$.

Before the end of 1984 there were two other major developments:

The heterotic string - Gross, Harvey, Martinec, Rohm

Calabi-Yau compactification - Candelas, Horowitz, Strominger, & Witten.

By the beginning of 1985, string theory — with the goal of unification — had become a mainstream activity.

THE END (& the beginning)